

Establishing Exciton Tunability in Two-Dimensional Perovskites

Kameron Hansen, October 2021

Introduction

With the advent of epitaxial deposition techniques, quantum well structures based on III-V semiconductors such were first fabricated in 1970s. Strong confinement effects within the quantum wells allowed for stable excitons, i.e. bound electron-hole pairs that are responsible for light absorption, emission, and charge transport and therefore are at the center of the quantum well's optoelectronic properties. The properties of stable excitons in these quantum wells were investigated in the 1980s and 1990s, leading to a Nobel prize and the development of LEDs, electro-optic modulators, lasers, and photodetectors which are still considered state-of-the-art.^{1,2}

Today, research efforts are focused on a new class of quantum well materials with stable excitons, namely two-dimensional metal halide perovskites (2D MHPs), which may be superior to III-V quantum wells for a variety of light-harvesting and light-emitting applications.³⁻⁶ Unlike III-V quantum wells, 2D MHPs naturally crystalize into their A_2BX_4 structure (shown in **Fig. 1**). This fundamental difference in the method of crystal growth has both advantages and disadvantages. On the one hand, the solution-based growth of 2D MHP crystals near room temperature is far more cost-effective and eco-friendlier than expensive epitaxial techniques. On the other hand, it was precisely the high level of control and tunability offered by epitaxial deposition which enabled the development of III-V quantum well technologies. Indeed, the success of these technologies is greatly attributed to advances in epitaxial methods that allowed for high tunability of the quantum well layer thicknesses, composition, and orientation, and crystalline quality.⁷⁻⁹ Researchers leveraged these variables to tune the exciton properties and create optimized devices.

While this pathway towards developing exciton tunability is unavailable for 2D MHPs, nature has instead offered an alternative route through the wide catalog of molecules and atoms that naturally crystalize in the A_2BX_4 structure. Indeed, over 50 different cations have successfully been incorporated into the A-site.⁶ The B-site metal atom can vary between lead, tin, and germanium while the X-site halide is commonly occupied by iodide, bromide, or chloride. While a wide range of chemical compounds have been synthesized and incorporated into device structures, there is a lack of understanding as to how chemical substitutions within the A_2BX_4 formula affect the material's exciton properties such as the binding energy, mass, diffusion length, lifetime, radius, dipole moment, and polarizability. Unfortunately, the same lattice effects (e.g. dynamic disorder and phonon-coupling) which give excitons in 2D MHPs their unique character and defect-tolerant properties also result in spectral structures which are difficult to interpret, and as a result, measurements of the exciton's properties are challenging.¹⁰ Nowhere is this more evident than with the difficulty in measuring exciton binding energies (E_B), a parameter of great importance for solar cell engineering as it determines the ratio of excitons to free carriers and for LED engineering as it strongly dictates electroluminescence quantum efficiency.¹¹ For example, taking the well-known 2D MHPs phenethylammonium lead iodide (PEA_2PbI_4) and butylammonium lead iodide (BA_2PbI_4) as case studies, state-of-the-art techniques for measuring E_B in 2D MHPs, namely absorption, photoluminescence excitation, electroabsorption, and low- and high-field magnetoabsorption at liquid helium temperatures have thus far led to E_B values ranging from 190 to 490 meV.^{12, 13} While some of this variance is caused by differences in the organic molecule, sample thickness, dielectric environment, and morphology, a significant portion of the variance originates from differing interpretations of band-edge absorption features and assumptions that are intrinsic to the measurement techniques. For theoretical and applied progress to continue forward, it is critical develop measurement techniques that are accurate and reproducible in determining the exciton's properties. Until the measurement uncertainty of E_B is less than the variance of E_B between different compositions, exploring compositional effects on E_B will be a fundamentally hopeless endeavor. Thus, it is expedient to develop novel measurement methods that are accurate, precise, and reproduceable.

In order to address these challenges, I propose to develop a highly accurate and precise method using electroabsorption (EA) spectroscopy to measure E_B in 2D MHPs. This new measurement method will then be applied to a wide range of 2D MHP compounds using electroabsorption (EA) and dielectric spectroscopies in order to measure the exciton's binding energy, radius, mass, dipole moment, and polarizability and thereby establish exciton tunability in 2D MHPs.

Specific Aims

1. Simulate the EA spectrum of a 2D Wannier exciton in order to better understand features in the EA spectrum of 2D MHPs and thereby develop a method for accurately measuring E_B in 2D MHPs.
2. Synthesize multiple 2D MHP compositions (elemental substitutions of A, B, and X into the generalized formula A_2BX_4) and use EA spectroscopy to measure the band gap (E_G) and exciton binding energy (E_B) for each composition.
3. Use dielectric spectroscopy to measure the real part of the dielectric function $\epsilon_r(\omega)$ for the barrier layer $[A_2]^{2+}$, well layer $[BX_4]^{2-}$, and the full 2D MHP material A_2BX_4 .
4. Use mathematical models, based on the image charge and quantum-confinement effects, to convert the E_B and ϵ_r measurements for 2D MHPs into theoretical values for 3D MHPs. Such models will be used to determine hypothetical E_B values for exfoliated 2D MHP sheets in a vacuum, and thereby contextualize our results within the universal scaling law for 2D systems, i.e. $E_B \sim E_G/4$.¹⁴

Two-dimensional metal-halide perovskites (2D MHPs) have become the focus of intense research efforts due to their unique charge-lattice interactions and impressive device performance in solar cells, LEDs, field-effect transistors, photodetectors, and polariton lasers.^{3, 15-18} For theoretical and applied progress to continue forward, it is critical for the community to reach a consensus of basic photophysical properties such as E_B and E_G ; however, state-of-the-art techniques to measure such values yield inconsistent results.^{12, 19} For my PhD research I propose to use EA and dielectric spectroscopies along with theoretical models for excitons in layered systems, as represented in **Fig. 1**, to not only obtain unambiguous measurements of these values but also to develop models that can be used to predict E_B in 3D MHPs. In accomplishing the above aims, the exciton binding energy will turn from an unknown into a controllable degree of freedom for future MHP device engineering.

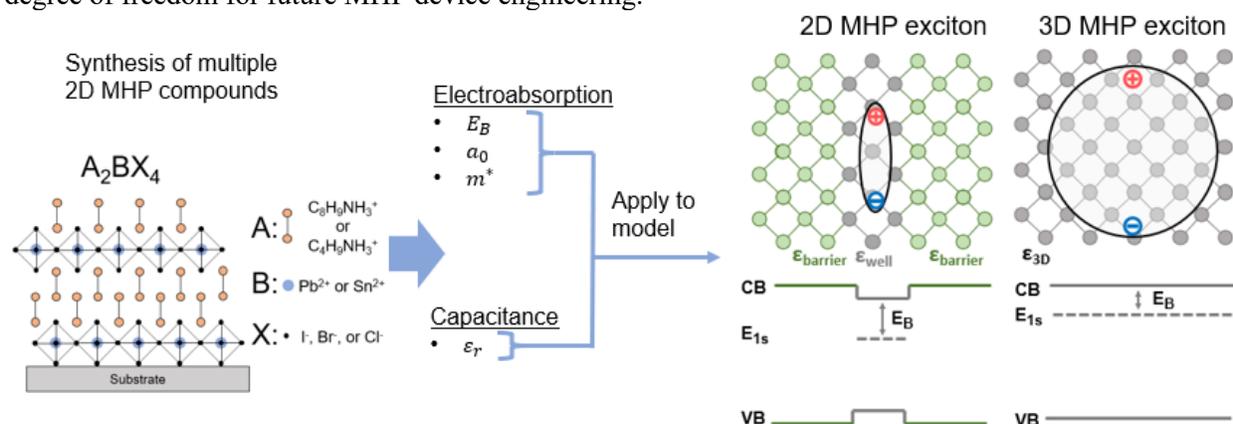


Figure 1: Pathway towards understanding the effect of chemical substitutions on the 2D MHP exciton properties. First, six compounds: PEA_2PbI_4 , PEA_2SnI_4 , PEA_2PbBr_4 , PEA_2PbCl_4 , PEA_2SnBr_4 , PEA_2SnCl_4 (where $PEA = C_8H_9NH_3^+$) will be synthesized. The measured EA spectrum will yield the exciton's binding energy (E_B), radius (a_0), and mass (m^*) while the measured capacitance will yield the material's dielectric constant (ϵ_r). These values will then be used to develop empirical models for E_B in 2D and 3D MHPs.

Research Plan

Aim 1: Numerical simulation of the 2D MHP EA spectrum

In pursuit of high confidence measurements of excitonic properties in 2D MHPs, I am motivated to turn to the technique that was most significant in advancing the fundamental understanding of III-IV quantum wells, namely electroabsorption spectroscopy (EA). EA measures the difference in a material's absorption spectrum, A , with and without an applied field:

$$\Delta A = A(F) - A(F = 0) \quad (1)$$

In the 1980's, the seminal work of Bastard, Mendez, and Miller,²⁰⁻²⁴ based on the theory of Dow and Blossey,^{25, 26} not only demonstrated the anisotropy of the inorganic quantum well electronic structure but also achieved a high level of agreement between theory and experiment²⁷ which lead to measurements of the exciton's Bohr radius a_0 , reduced effective mass m^* , and E_B .

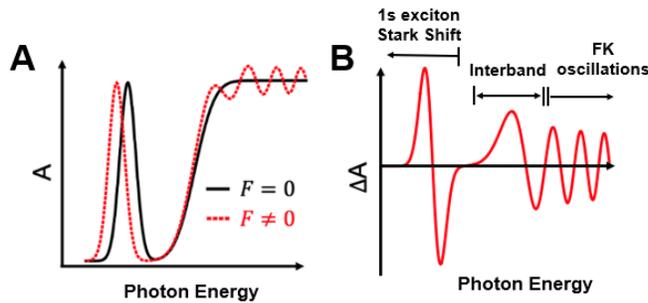


Figure 2: (A) Prototypical absorption (black) and field-shifted absorption of an excitonic semiconductor and the corresponding (B) EA signal $A(F = 0) - A(F \neq 0)$.

The standard absorption spectrum (Fig. 2(A)) and EA spectrum (Fig. 2(B)) for an excitonic semiconductor is divided into three regions: (1) below-gap where the exciton states redshift according to the Stark effect (2) near-gap where interband transitions are affected by the Franz-Keldysh (FK) effect and (3) above-gap where the same FK effect produces sinusoidal-like oscillations over a wide spectral range. The EA response of 2D MHPs closely resembles this behavior.²⁸⁻³¹ While the Stark shift of the 1s exciton is unambiguous in all reports of 2D MHP EA spectra, the near-gap and above-gap EA features are less straightforward making it

difficult to pinpoint E_G and leading to contradictory measurements of E_B (190 vs 361 meV for PEA_2PbI_4 ^{28, 29} and 220 vs 364 meV for BA_2PbI_4 ^{30, 31}). It is important to note that the variance here originates from differing interpretations of nearly identical EA spectra. Indeed, the extremely close qualitative resemblance of the 2D MHP EA response in these four studies demonstrates the measurement is reliable and it is a lack of understanding, not measurement error, that has thus far impeded EA from producing reliable measurements of the exciton's properties. Therefore, the first aim of this study is the bridge the gap between the experimental EA features and our theoretical understanding of them by simulating the EA response of a 2D Wannier exciton in a uniform electric field in a regime that should be expected for 2D MHPs, i.e. a low-field and low-broadening regime. The same theory, in a high-field low-broadening regime was extremely successful in describing the EA response of III-V quantum wells.²⁷ However, this theory has thus far been overlooked in its application to 2D MHPs, perhaps in-part due to the fact a lengthy and complex numerical procedure is needed to solve the Schrödinger equation of a 2D Wannier exciton in a uniform field. In dimensionless units, this equation takes the following form:²⁶

$$\left(-\nabla_{x,y}^2 - \frac{2}{r} + fx\right)\psi_n(x,y) = E_n\psi_n(x,y) \quad (2)$$

where $\psi_n(x,y)$ is the wave function for the electron-hole relative motion, E_n is $\hbar\omega - E_G$, f is a dimensionless field defined as the ratio between the applied electric field and the field required to create a potential drop of 1 E_B across the exciton's radius (i.e. the ionization field), and is defined in terms of the material parameters, as follows:²⁵

$$f = \frac{e^5}{128 \pi^3 \hbar \epsilon_0} \frac{(m^*)^2}{(\epsilon^*)^3} \quad (3)$$

where m^* is the exciton's reduced effective mass, ϵ^* is the effective dielectric constant,³² and the physical constants e , \hbar , and ϵ_0 have their standard definitions. Assuming reasonable values for ϵ^* between 3 – 5 and a reduced effective mass m^* between 0.050 – 0.221 m_0 ,^{12,33} the ionization field for 2D MHPs is calculated to fall the range of $10^2 - 10^4$ kV/cm, far greater than the maximum field strength that can be produced experimentally (~ 30 kV/cm). Therefore, the 2D MHP EA spectrum should be simulated in the low-field regime where ionization effects are negligible $0.003 < f < 0.3$. The final material parameter needed for the simulation is the homogenous broadening Γ which I have measured to be 11 meV by fitting a Lorentzian profile to the 1s exciton absorption peak.

To make **Eq. (2)** tractable, a transformation to parabolic coordinates is needed. This results in two quasi-Schrödinger equations with boundary conditions that have been determined by Dow and Blossy.²⁵ ²⁶ My proficiency in using Matlab to solve differential equations will allow me to implement the numerical procedure that is described in detail in **Ref. (34)** and **Ref. (35)** in order to determine the electron-hole wavefunctions $\psi_n(r)$ for a range of photon energies E and field strengths f .^{34, 35} These wavefunctions relate to the semiconductor's absorption spectrum $\alpha(E)$ through Elliot's formula for direct interband transitions:^{36, 37}

$$\alpha(E, f) = B \sum_n |\psi_n(r = 0, f)|^2 \delta(E_n - \hbar\omega) \quad (4)$$

where B is a constant based on material parameters that can be set to unity for the purposes of simulating the EA line shape. The simulated EA spectrum is then calculated as:

$$\Delta\alpha = \alpha(E, f) * L(E) - \alpha(E) * L(E) \quad (5)$$

where $L(E)$ is a Lorentzian function with a half-width-half-maximum (HWHM) of Γ . The resulting EA features of the 2D Wannier exciton and their positions, relative to E_G and E_{1s} , will allow for precise measurements of E_B and a deepened understanding of the 2D MHP EA response.

Aim 2: Measure the exciton properties for multiple A_2BX_4 compositions

The results of the numerical simulation will be compared to the experimental EA spectra for a wide range of 2D MHP chemical compounds with general formula A_2BX_4 . In the last decade, many photophysical studies on the prototypical 2D MHP material PEA_2PbI_4 have led to a deepened understanding of the exciton physics and dynamic coupling to phonon modes within the $[PbI_4]^{2-}$ layer.^{10,38} Simultaneously, many researchers have turned to the tin-based PEA_2SnI_4 for the fabrication of field effect transistors (FET) and bromide and chloride varieties of PEA_2PbBr_4 and PEA_2PbCl_4 for the fabrication of UV photodetectors and LEDs.^{4,17,39,40} It is well established that the band gap increases moving from I to Br to Cl, and decreases moving from Pb to Sn. However, outside of the band gap, the effect of these chemical substitutions on the electronic structure and exciton properties remains largely unexplored, particularly from an experimental perspective.

First principle calculations starting with Hartree-Fock studies in the 1990s have been in agreement the occupant of the B-site metal center is most consequential in determining electronic structure.⁴¹⁻⁴³ The conduction band primarily consists of the B-site metal's highest-occupied p orbital, meanwhile the valence bands are largely comprised of the X-site halide's highest-occupied p orbital anti-bonded to the B-site metal's highest-occupied s orbital. On the other hand, the A-site organic molecule offers a valuable degree of tunability by modulating the dielectric environment and phonon modes, but has minimal effect on the band structure.^{10, 44, 45} Therefore, to efficiently explore the space of electronic structure and exciton tunability in 2D MHPs, I will synthesize the following six compounds: PEA_2PbI_4 , PEA_2SnI_4 , PEA_2PbBr_4 , PEA_2PbCl_4 , PEA_2SnBr_4 , PEA_2SnCl_4 . I choose phenethylammonium (PEA) as the A-site cation not only

because it is the most commonly studied, but also due to the fact 2D MHP crystals with PEA maintain a single structural phase between room and liquid-helium temperatures, which is rare for MHP systems.⁴⁶ This lack of phase transition will be ideal for EA measurements as it allows the absorption broadening to be reduced at low temperatures without introducing new spectral structures, e.g. additional exciton peaks from secondary crystalline phases. Wherein the EA spectra of these six compounds fails to produce a clear trend, supplementary compounds will be synthesized using a butylammonium (BA) cation in place of PEA. Since the electronic coupling between the molecule and semiconducting layers is negligible, comparison of BA-based compounds will serve as a secondary data point in elucidating the effect of B-site and X-site cations on the excitonic properties.

Each of the 2D MHP compositions will be synthesized using established spin-coating techniques to produce polycrystalline thin films.^{47,48} In brief, the ammonia salt PEAX or BAX (where X = I, Cl, or Br) will be combined with a corresponding lead or tin salt with a matching halide (PbX_2 or SnX_2), dissolved in a 4:1 DMF:DMSO solvent, and pipetted onto a substrate for spin coating. It has been shown that this solvent ratio produces 2D MHPs with quantum well layers homogeneously stacked parallel to the substrate.⁴⁹ Achieving this orientation is critically important for probing the exciton's in-plane EA response, and therefore, my colleagues and I will collect grazing-incidence X-ray diffraction (GIXD) patterns in order to measure the distribution of crystallite orientations within the polycrystalline film.⁵⁰ The EA signal of the 2D MHP films will also be affected by the film's microstructure and average grain size. Disordered films will have broader absorption features, and therefore, smaller EA signals. Furthermore, small grains will reduce the probability of observing FK oscillations above the band gap, which are needed to measure needed for measuring m^* . These oscillations persist for a spectral range that is proportional to the mean free path of conduction band electrons, and therefore, will be absent in films with small grain sizes (< 100 nm for MHPs).⁵¹ Thus, in order to target large grains and continuous-film microstructure for each composition, I will test three common spin-coating procedures (hotcast⁴⁸, antisolvent,⁵² and post-annealed⁵³) and characterize the resulting film quality using optical microscopy, atomic force microscopy (AFM), and X-ray diffraction. The procedure that yields the highest quality film will then be used to spin-coat films for EA and dielectric spectroscopies.

The proposed methods for measuring the 2D MHP EA and dielectric spectra are highlighted in **Fig. 3**. An array of interdigitated fingers (**Fig. 3(A)**) will be used to produce high electric fields (5 – 30 kV/cm) parallel to the substrate while also allowing for the transmission of light. Postdoctoral scholar Hao-Chieh Hsieh will use photolithography to deposit gold electrodes onto quartz substrates in the University of Utah clean room. Perovskite films will then be spin-cast directly onto the substrate array for EA (**Fig. 3(B)**) and dielectric measurements (**Fig. 3(C)**). The EA procedure will be carried out in Dr. John Colton's laboratory in Brigham Young University and is as follows: first, samples are mounted into a cryostat, cooled to low temperature in order to reduce the homogenous broadening, and then the transmission of sample T , transmission of substrate T_0 , and electro-transmission of sample ΔT are collected as independent wavelength scans (set-up depicted in **Fig. 3(B)** for ΔT mode). The absorption $A(E)$ and EA spectra $\Delta A(E)$ will then be calculated as:⁵⁴

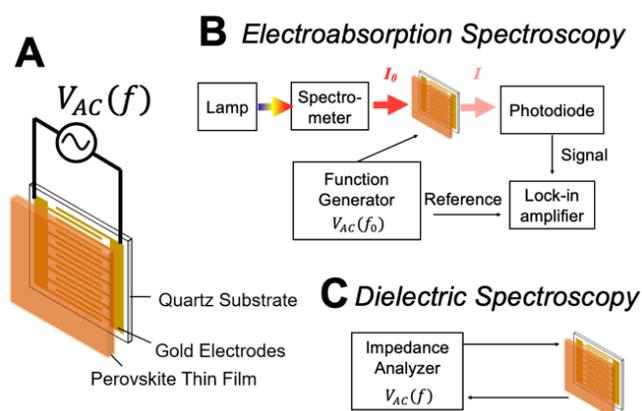


Figure 3: (A) Perovskite film cast on a substrate of interdigitated electrodes which can be used to measure the film's (B) EA response and (C) dielectric spectrum.

$$A = \log_{10}(T_0/T) \quad (6)$$

$$\Delta A = -\log_{10} \left(1 + \frac{\Delta T}{T} \right) \quad (7)$$

The measured EA spectrum $\Delta A(E)$ will be compared to the simulated EA spectra $\Delta\alpha(E)$ for each 2D MHP composition. The relative alignment critical points in $\Delta A(E)$, relative to E_{1s} and E_G in $\Delta\alpha(E)$, will allow for determination of $E_B = E_G - E_{1s}$. Once E_B is measured with high-confidence, the field dependence of the Stark shift can be used to measure the exciton's transition dipole moment μ , polarizability α , and Bohr radius a_0 .⁵⁵ If high-field Franz-Keldysh features are resolved, the exciton's reduced effective mass m^* can also be determined.⁵⁶ However, measurement of a_0 , μ , α , and m^* requires knowledge of the electric field strength F , as experienced by the exciton, which depends on the real part of the dielectric function ε_r at the frequency of the modulating voltage signal, $V_{AC}(f_0)$. The dielectric properties of 2D MHPs have thus far not been reported in the literature, and therefore, it is necessary to develop a method to measure $\varepsilon_r(f_0)$ for each 2D MHP thin film.

Aim 3: Measure dielectric values for multiple A_2BX_4 compositions

Measuring the dielectric spectrum for each 2D MHP composition serves two critical roles in accomplishing the aims of this proposed research. First, knowledge of $\varepsilon_r(f_0)$ will allow the exciton's properties a_0 , μ , α , and m^* to be measured from the EA spectrum, as previously discussed. Second and equally importantly, the dielectric value near 10^7 Hz is an excellent measurement of the effective dielectric value ($\varepsilon_r(10^7 \text{ Hz}) \approx \varepsilon^*$)^{32, 57} which is an important input parameter for the mathematical models of **Aim 4**. This is because E_B is primarily determined by ε^* and m^* , with the former carrying greater weight.¹⁴

Traditionally, the dielectric properties of MHPs are measured by pasting silver electrodes onto single crystals and recording the capacitance of the crystal in a parallel-plate geometry for a spectrum of AC frequencies. However, as we show in the following, the substrates with interdigitated electrode fingers can also be employed to directly measure dielectric properties of 2D MHP thin films (represented in **Fig. 3(C)**). This geometry will allow for facile dielectric measurements of the long list of thin film materials needed for this study, namely: PbX_2 , SnX_2 , PEA_2PbX_4 , PEA_2SnX_4 , $\text{C}_3\text{NH}_3\text{PbX}_3$, $\text{C}_3\text{NH}_3\text{SnX}_3$ where $X = \text{I, Br, or Cl}$ (18 films total), without the need of growing single crystals. Thin films of metal-halide salts (PbX_2 , SnX_2) and 3D MHPs ($\text{C}_3\text{NH}_3\text{PbX}_3$, $\text{C}_3\text{NH}_3\text{SnX}_3$) can be created with spin-coating procedures similar to the 2D MHP procedure,^{48, 58} and the measured dielectric values will be used as input parameters for the mathematical models of **Aim 4**. As first shown by den Otter,⁵⁹ the capacitance (C) of interdigitated electrodes in a uniform dielectric medium can be well-approximated by a summation of zeroth Bessel function of the first kind (J_0). My collaborator Emma McClure modified den Otter's equations to fit a thin film geometry in order to account for the independent dielectric loads of the quartz substrate ($\varepsilon_{r,q} = 3.7$),⁶⁰ the thin film ($\varepsilon_{r,f}$), and air ($\varepsilon_{r,a} = 1$), as follows:

$$C = \frac{pq\pi}{8V_0a} \sum_{n=1}^{\infty} \frac{J_0\left(\frac{(2n-1)\pi s}{2a}\right)}{2n-1} \left[\varepsilon_{r,q}(1 - e^{-(2-4n)\pi t_q/a}) + \varepsilon_{r,f}(1 - e^{-(2-4n)\pi t_f/a}) + \varepsilon_{r,a}(e^{-(2-4n)\pi t_q/a} + e^{-(2-4n)\pi t_f/a}) \right] \quad (8)$$

where p , q , s , and a denote various dimensions of the interdigitated electrodes and t_q and t_f denote the thickness of the quartz substrate and thin film, respectively. Optical profilometry and AFM in the University of Utah Nanofab will be used to measure t_f . The 2D MHP thin films will be cast onto interdigitated electrodes and electrically connected to the input of an impedance analyzer. The capacitance across a wide frequency range ($10^0 - 10^7$ Hz) will be recorded, and with the use **Eq. (8)**, these capacitance values will be used to obtain $\varepsilon_{r,f}$.

Aim 4: Mathematical model for E_B in MHPs

While the hydrogenic formalism of **Eq. (2)** represents the state-of-the-art theoretical description of a 2D MHP exciton in an electric field, more sophisticated theoretical models exist for describing 2D MHP excitons in the absence of external electric fields.^{12, 61, 62} Excitons in 2D MHPs deviate from the hydrogen model primarily due to: (1) their non-zero spatial extent in the \hat{z} direction which results in a level of quantum confinement that is less than the 2D case and (2) their anisotropic dielectric environment which results in an E_B -enhancing image charge effect. In a simple model, the hydrogenic expression for E_B can be modified to account for these two effects, as follows:

$$E_B = \left(\frac{\epsilon_w}{\epsilon_b}\right)^\gamma \left(\frac{2}{\alpha - 1}\right)^2 \left(\frac{1}{\epsilon^*}\right)^2 \left(\frac{m^*}{m_0}\right) R_H \quad (9)$$

where ϵ_w is the dielectric of the well layer ($[\text{BX}_4]^{2-}$), ϵ_b is the dielectric of the barrier layer ($[\text{PEA}_2]^{2+}$), γ is a real number that has yet to be empirically determined, α is the fractional dimensionality, m_0 is the electron mass, and R_H is the Rydberg energy of hydrogen. The first factor $(\epsilon_w/\epsilon_b)^\gamma$ in **Eq. (9)** accounts for the image-charge effect on E_B ^{61, 63} while the second factor $(2/(\alpha - 1))^2$ accounts for quantum confinement in a fractional-dimensional (α -D) system.⁶⁴ In a purely 2D system, the dimensionality $\alpha = 2$ results in a factor of 4 enhancement and $E_{B,2D} = 4E_{B,3D}$. In 2D MHPs, however, the nonzero thickness of the well layer leads to $\alpha = 2.3$, as measured by Tanaka et al.⁶⁵ The values of E_B and m^* will be known from the EA spectrum, ϵ_w will be approximated by the dielectric constants of the PbX_2 and SnX_2 thin films, ϵ_r (10^7 Hz) of the 2D MHP film will be taken as ϵ^* , and the value $\epsilon_b = 3.32$ is known from dielectric measurements of PEA in the literature.³³ Therefore, all variables in **Eq. (9)** will be known, except for γ , for six 2D MHP compositions. Thus, γ will be empirically determined and the validity of **Eq. (9)** will be established by the standard deviation of γ across the six 2D MHP compositions.

Provided that **Eq. (9)** is found to be a good approximation, the E_B values for 2D MHPs can be converted to theoretical E_B values for 3D MHPs, as follows:

$$E_{B,2D} = \left(\frac{\epsilon_w}{\epsilon_b}\right)^\gamma \left(\frac{2}{\alpha - 1}\right)^2 \left(\frac{\epsilon_{3D}^*}{\epsilon_{2D}^*}\right)^2 E_{B,3D} \quad (10)$$

where ϵ_{3D}^* will be known from dielectric measurements of $\text{CH}_3\text{NH}_3\text{PbX}_3$ and $\text{CH}_3\text{NH}_3\text{SnX}_3$. This expression assumes $m_{2D}^* = m_{3D}^*$ which has been found to be a good approximation by numerous computation and experimental studies^{12, 42, 66-68} due to the fact the 2D and 3D MHP electronic structures originate from the same metal-halide orbitals. Thus, the tunability of E_B for the class of 3D MHP materials can be reasonably estimated from E_B measurements on 2D MHPs. This is important because, while 3D MHPs are extremely impactful optoelectronic materials,⁶⁹ the variance in experimental and computational estimates for E_B ($5 - 50$ meV)⁷⁰⁻⁷² far exceeds the variance between different chemical compositions.

In addition to the calculations detailed above, I will also reach out to establish a collaboration with Dr. Jacky Even and Dr. Claudine Katan who are the world's expert in modeling dielectric and quantum confinement in 2D MHP using state-of-the-art theoretical models. The E_B , a_0 , and m^* from the EA spectrum can be used to improve Dr. Even's and Dr. Katan's models of the dielectric anisotropy in 2D MHPs. These improved models can then be used to predict the magnitude of E_B enhancement expected for exfoliated single layers on a variety of substrates and no substrate ($\epsilon_b = 1$).

Proposed Timeline:

Aim 1	Develop simulation	
Aim 2	Optimize fabrication	Collect EA
Aim 3	Spin-coat 2D and 3D MHPs	Measure ϵ_r
Aim 4	Identify state-of-the-art models	Calculate dielectric effects on E_B
Collaboration		
	Year 1	Year 2

Challenges and Alternative Approaches

Experimental artifacts in EA spectroscopy

In general, modulation spectroscopies such as EA are trustworthy since the normalized difference signal $\Delta T/T$ is free of the instrument response function. However, modulation of the transmission can sometimes result from phenomena other than field-induced changes in the absorption. These phenomena include: (1) field-induced changes in the refractive index (Δn) and (2) thermal-induced changes in the absorption from Joule heating. I plan to calculate the magnitude of these two effects in order to estimate their contribution to the EA response of 2D MHPs. A Taylor expansion of the EA signal in terms of small changes in the refractive index Δn and absorption coefficient $\Delta\alpha$ results in the following expression:⁷³

$$\Delta A = C_1\Delta n + C_2\Delta\alpha - \Delta\alpha t_f \quad (11)$$

where C_1 and C_2 are optical constants that depend on the real n and imaginary κ components of the refractive index (expressions given in **Ref. (73)**) and t_f is the film thickness. Thus, the influence of field-induced changes in the 2D MHP refractive index on the EA signal can be estimated provided that n and κ are known across the spectral range of interest. I plan use a Kramer's Kroning analysis of the transmission data to calculate the n and κ , according to the method in **Ref. (74)**.⁷⁴ The values of C_1 and C_2 will then be computed to determine their contribution to the EA signal. If refractive index effects are found to be large, I will consider using a simulated Δn (by a Kramer's Kroning transformation of the simulated $\Delta\alpha$) to subtract $C_1\Delta n$ from the EA signal.

The contribution of thermal effects to the EA signal can be estimated by:⁷³

$$\Delta A_{thermal} = \frac{dA}{dE} \frac{dE_G}{dT} \Delta T \quad (12)$$

where T now denotes temperature, not transmission. ΔT is the RMS modulation of the temperature at the lock-in's reference frequency f_0 and dE_G/dT is the rate of the band gap's blueshift with temperature (~ 0.2 meV/K for 2D MHPs).⁷⁵ The value of ΔT can be calculated from a power-balance equation: Power in – Power out = CdT/dt where C is the heat capacity of the substrate. The ‘power in’ terms include contributions from light absorption and Joule heating, the latter of which will be larger for the tin-based perovskites due to their higher electrical conductivity, while the ‘power out’ will represent the thermal contact with the cold finger. I will set up this power balance equation, solve for $T(t)$, and use the RMS value at the modulation frequency (ΔT) to calculate $\Delta A_{thermal}$ in **Eq. (12)**. If thermal effects are found to be large, I will consider using chemical additives to increase the resistivity of the perovskite film.⁷⁶

High measurement uncertainty

The proposed method for measuring m^* requires that FK features in the EA spectrum show a clear and consistent rate of field-broadening,⁵⁶ which may not be the case for all 2D MHP compositions. If the clarity and reproducibility of field-broadening within the FK features is poor, then I will establish a collaboration with an expert in using first principle approaches, such as density functional theory (DFT), to model MHP systems. DFT has proven effective at accurately determining m^* in MHP systems from the band curvature at $k = 0$.⁷⁷

Another measurement that may prove to have high uncertainty is ϵ_w , i.e. the dielectric value of the well layer ($[BX_4]^{2-}$). Using the $\epsilon_{r,f}(10^7 \text{ Hz})$ value of a BX_2 thin film to estimate ϵ_w may be an unfriendly approximation due to differing LO phonon contributions to ϵ_r between BX_4 2D sheets and BX_2 3D lattices. Fortunately, the validity of this measurement can be easily assessed using the following expression for the quantum well's dielectric value:⁶¹

$$\epsilon_r = \frac{\epsilon_w L_w + \epsilon_b L_b}{L_w + L_b} \quad (13)$$

where L_w and L_b are the length of the well and barrier layers, respectively. If ϵ_r from Eq. (13) differs from the measured dielectric value for multiple of the six 2D MHP compositions, then it will be known that using BX_2 films to estimate ϵ_w is a poor approximation. In such a case, ϵ_w will also be determined using DFT modeling. While the magnitude of the polarization contributions from valence electrons and phonons to a material's dielectric constant is sometimes difficult to accurately determine with first principle approaches, DFT calculations are comparatively more effective at capturing relative differences between materials, provided the same methodology is used. Therefore, ϵ_b for the barrier layer $[PEA_2]^{2+}$ will be calculated along with ϵ_w to arrive at an accurate ratio (ϵ_w/ϵ_b) that is needed for Eq. (9) and Eq. (10).

Preliminary Results:

My colleagues and I have made significant progress carrying out the proposed research for the commonly-researched PEA_2PbI_4 material. In Fig. 3, the measured absorption and EA spectra of PEA_2PbI_4 (Fig. 3(A)) is shown in comparison the theoretical absorption and EA spectrum (Fig. 3(B)) simulated using the theory of Dow and Lederman.³⁴ The measured and simulated spectra have a high level of qualitative agreement. It was found that the zero-crossing of the exciton and interband (IB) features, labeled at 'b' and 'e' respectively in Fig. (3A), are the correct markers of E_{1s} and E_G . Thus, $E_B = E_G - E_{1s}$ (Δbe) was measured to be 223 ± 3 meV. Turning to reported EA spectra in the literature, we find that the positions of 'b' and 'e' are remarkable consistent across various studies which employ different procedures for sample fabrication and for EA measurements. As summarized in Table 1, our Δbe measurement method yields E_B values with 1σ variance of ~ 3 meV for these EA spectra of PEA_2PbI_4 and BA_2PbI_4 in the literature, thus demonstrating that EA is a remarkable effective and reproducible technique when applied to 2D MHPs.

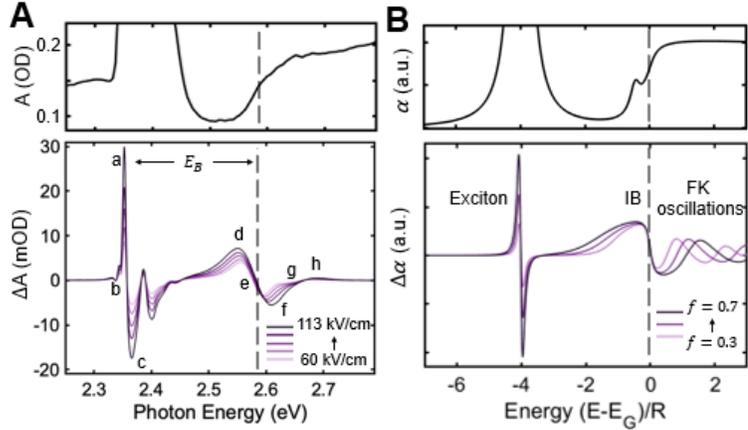


Figure 4: (A) Absorption (top) and EA (bottom) spectra for PEA_2PbI_4 at 14 K in comparison with (B) the simulated absorption (top) and EA (bottom) spectra for a 2D Wannier exciton in the low-field, low-broadening regime.

Table 1: Comparison of E_B values from the literature, as measured by the Δbe method

Material	Temp. (K)	E_{1s} (b) (meV)	E_G (e) (meV)	E_B (Δbe) (eV)	Reported E_B (eV)	Ref
PEA_2PbI_4	15	2.357	2.579	223 ± 3	223	Hansen et al ⁵⁴
PEA_2PbI_4	50	2.38	2.60	220 ± 5	190	²⁸
BA_2PbI_4	45	2.551	2.802	251 ± 2	251	Hansen et al ⁵⁴
BA_2PbI_4	50	2.550	2.800	250 ± 5	220	³⁰
BA_2PbI_4	45	2.561	2.814	253 ± 5	340	³¹

The dielectric value of PEA_2PbI_4 was measured using both parallel plate and interdigitated electrode geometries. The two methods were in agreement and we found $\epsilon_r(f_0)$ to be 4.5. Knowledge of E_B and $\epsilon_r(f_0)$ allowed for determination of the exciton's Bohr radius ($a_0 = 2.2 \pm 0.55$ nm), transition dipole moment ($\mu = 13 \pm 4$ D), and polarizability ($\alpha = 80000 \pm 10300$ Å³) according to the field dependence of the exciton's Stark shift as well as the reduced effective mass ($m^* = 0.09 \pm 0.024 m_0$) according to the field

dependence of IB feature. These results have been summarized in a manuscript which is under review at *PRX Energy* and currently available for viewing on arXiv.org.⁵⁴

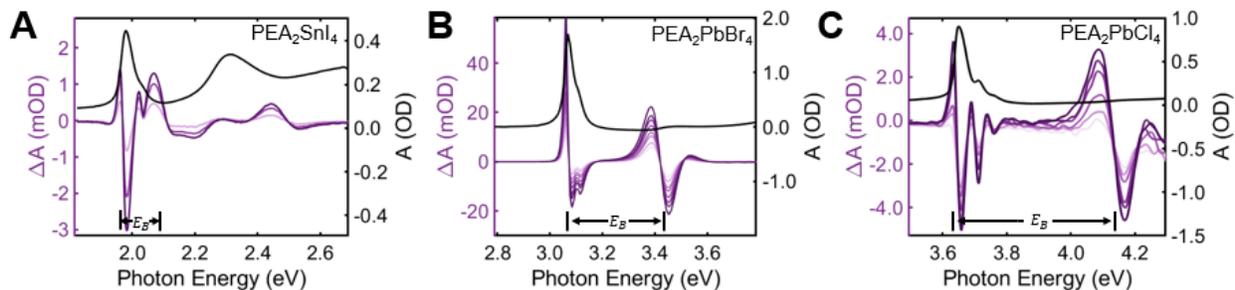


Figure 5: EA (purple) and absorption (black) spectra for (A) PEA_2SnI_4 showing $E_B = 100$ meV, (B) $\text{PEA}_2\text{PbBr}_4$ showing $E_B = 348$ meV, and (C) $\text{PEA}_2\text{PbCl}_4$ showing $E_B = 487$ meV.

While the dielectric constants of other 2D MHP compositions have yet to be measured, I have measured low-temperature EA spectra for three other compositions: PEA_2SnI_4 (**Fig. 5(A)**), $\text{PEA}_2\text{PbBr}_4$ (**Fig. 5(B)**), and $\text{PEA}_2\text{PbCl}_4$ (**Fig. 5(C)**) resulting in accurate and precise determination of E_G and E_B . In **Fig. 6**, I plot relationship between E_G and E_B for 2D MHPs in comparison to the universal scaling law for 2D materials, that is, $E_B \sim E_G/4$.¹⁴ It is clear the E_B values in 2D MHPs are highly tunable, from 100 meV for PEA_2SnI_4 to 487 meV for $\text{PEA}_2\text{PbCl}_4$ and that 2D MHPs in bulk form give access to a range of parameter space (low E_G , high E_B) that is unique for semiconductors. This is undoubtedly due to the factor of ~ 3 enhancement to E_B that occurs as a result of the image charge effect. These data serve as unambiguous correction to previous studies that report E_B decreases moving from I to Br to Cl and stays relatively constant moving from Pb to Sn.^{19, 41, 63}

Upon obtaining dielectric measurements and EA measurements for a wider range of 2D MHP compositions, my colleagues and I will be able to empirically determine γ , i.e. the power relation between the dielectric mismatch (ϵ_w/ϵ_b) and the magnitude of the image charge effect on E_B . Using $\epsilon_b = 1$, it will be possible to predict E_B for exfoliated sheets of 2D MHPs in a vacuum and compare the resulting E_B values to the universal trendline $E_B \sim E_G/4$.

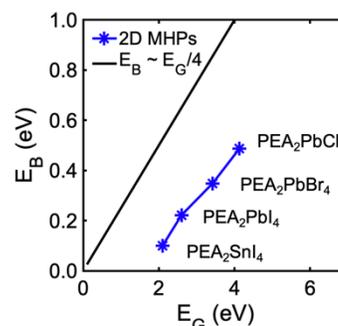


Figure 6: 2D MHPs in comparison to the universal scaling law for 2D materials $E_B \sim E_G/4$.

Conclusion/Outlook:

Two- and three-dimensional MHPs are the star optoelectronic materials of this decade, but the development of tunable optoelectronic properties in MHPs has thus far been limited. Unlike III-V quantum wells, compositional changes to the 2D MHP structure is the primary pathway available for tuning the exciton's properties in this class of quantum wells. However, the variance between reported E_B values for a given A_2BX_4 composition (in the 2D case) or a given ABX_3 composition (in the 3D case) is the same order of magnitude as the expected variance between different chemical compositions. Therefore, establishing exciton tunability in MHP systems is a fundamentally hopeless endeavor until a precise, accurate, and reproducible method for measuring E_B is developed.

The four aims of this proposed work directly address this challenge and establish exciton tunability in MHPs. I will simulate the 2D MHP EA spectrum in **Aim 1** to gain a deeper understanding of the features within the experimental EA spectrum of 2D MHPs. The preliminary results show that this leads to consistent and reliable measurements of E_B . In **Aim 2**, I will exploit the sensitivity of EA spectroscopy and the large spectroscopic separation of exciton and continuum states in 2D MHPs to obtain high-confidence

measurements of E_B for a wide range of A_2BX_4 compositions (A = PEA, B = Pb or Sn, X = I, Br, or Cl). These measurements will be complemented with dielectric spectroscopy in **Aim 3** and the measured values will be used at input parameters for empirical (and potentially theoretical) models, as described in **Aim 4**, to establish exciton tunability in this important class of materials. These are important steps towards the development of next-generation photovoltaics, LEDs, and other optoelectronic devices.

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